Simulation of Nondifferentiable Models for Groundwater Flow and Transport

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Outline

- Nonsmooth models
 - Richards' equation: van Genuchten/Mualem formulae
 - Reactive transport Freundlich isotherm
- What solvers must do
 - ODE/DAE formulations
 - Nonsmooth calculus and ADH
 - Temporal error estimation and control (time?)
- Conclusions

Collaborators

- ERDC: Stacy Howington, Charlie Berger, Jackie Hallberg
- NCSU: Jill Reese
- UNC: Casey Miller, Matthew Farthing, Joe Kanney
- Clemson: Lea Jenkins
- Mathworks: Mike Tocci
- Old Dominion: Glenn Williams

Richards' Equation: pressure head form

$$S_s S_a(\psi) \frac{\partial \psi}{\partial t} + \eta \frac{\partial S_a(\psi)}{\partial t} = \nabla \cdot [K(\psi) \nabla (z + \psi)]$$

		C	an a sifin atomana
ψ	pressure nead	\mathfrak{Z}_{S}	specific storage
$S_a(oldsymbol{\psi})$	saturation	η	porosity
$K(oldsymbol{\psi})$	hydraulic conductivity		

van Genuchten and Mualem formulae

$$S_a(\psi) = egin{cases} S_r + rac{(1-S_r)}{[1+(lpha|\psi|)^n]^m}, & \psi < 0\ 1, & \psi \ge 0 \end{cases},$$

$$K(\psi) = \begin{cases} K_s \frac{[1 - (\alpha |\psi|)^{n-1} [1 + (\alpha |\psi|)^n]^{-m}]^2}{[1 + (\alpha |\psi|)^n]^{m/2}}, & \psi < 0\\ K_s, & \psi \ge 0 \end{cases}$$

- S_r residual saturation
- α coefficient for mean pore size
- *K_s* saturated hydraulic conductivity
- *n* measure of pore size uniformity; m = 1 1/n

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- ERDC ADH code uses PL splines Lipschitz continuous/not differentiable

Reactive Transport in Porous Media

Freundlich isotherm:

$$\frac{C_s\eta}{\rho_b} = K\max(C,0)^r$$

Transport equation:

$$(C + \frac{\rho_b}{\eta} K \max(C, 0)^r)_t + \nabla \cdot [C \mathbf{v} - \mathbf{D} \nabla C] = 0$$

- C_s equilibrium concentration in the solid phase
- *r* Freundlich exponent
- C Freundlich coefficient
- ρ_b bulk density of the solid phase
- η porosity
- v mean pore velocity
- **D** hydrodynamic dispersion tensor

Nonsmoothness and a Fix

Nonlinearity is not Lipschitz continuous if 0 < r < 1. Fix: Differential Algebraic Equation (DAE) formulation Differential equation:

$$m_t + \nabla \cdot [C\mathbf{v} - \mathbf{D}\nabla C] = 0.$$

Algebraic constraint:

$$\left(\frac{\eta \max(m-C,0)}{\rho_b K}\right)^{1/r} - C = 0.$$

And now everything is differentiable,

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DAE and ODE Dynamics

$$u_t = f(t, u), \quad u(0) = u_0$$
 ODE

$$f(t, u, u_t) = 0, \quad u(0) = u_0, u'(0) = u'_0$$
 DAE

We have at most Index-one DAEs here. i.e. Implicit Euler works. Initial data for u'(0) is the solver's job.

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- Make temporal integration work better (Richards)
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It's still your job to design good solvers

- regularity of the solution
- differentiability of the nonlinearity
- discretizations
- linear solvers and preconditioning

DAE formulation of Reactive Transport Equation

Two equations for m and C

 $m_t = -\nabla \cdot [C\mathbf{v} - \mathbf{D}\nabla C] = f(m, C)$ Differential Equation

and

$$\left(\frac{\eta \max(m-C,0)}{\rho_b K}\right)^{1/r} - C = g(m,C) = 0$$
 Algebraic Constraint

There's no C_t anywhere.

Solving Reactive Transport Equation with Implicit Euler

Discretize in space, and advance in time by solving

$$m^{n+1} = m^n + hf(m^{n+1}, C^{n+1}),$$

$$g(m^{n+1}, C^{n+1}) = 0.$$

So the equation is for $u = (m, C)^T$.

Newton's method

Solve

$$F(u) = 0$$

by

$$u^+ = u^c + s, \qquad F'(u_c)s = -F(u_c)$$

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- Gaussian elimination (compute and factor matrix)
- iterative method with computed (approximate) Jacobian
- Matrix free: iterative method, finite difference Jacobian-vector products

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Everything works if $F'(u^*)$ is nonsingular.

What do you feed the solver?

$$F\left(\begin{array}{c}m\\C\end{array}\right) = \left(\begin{array}{c}m-m^n - hf(m,C)\\g(m,C)\end{array}\right) = \left(\begin{array}{c}0\\0\end{array}\right)$$

Solve with Newton. Converged result is $(m^{n+1}, C^{n+1})^T$.

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Solve with Newton. Converged result is $(m^{n+1}, C^{n+1})^T$. For small *h*,

$$F' = \left(\begin{array}{cc} I - hf_m & -hf_C \\ g_m & g_C \end{array}\right)$$

is nonsingular if g_C is nonsingular (aka index one).

Is RE a DAE?

Discretize in space, and you have

$$S_s S_a(\psi) \frac{\partial \psi}{\partial t} + \eta \frac{\partial S_a(\psi)}{\partial t} = N(\psi)$$

ODE solve: Use the chain rule and get

$$\frac{\partial \psi}{\partial t} = \frac{N(\psi)}{S_s S_a(\psi) + \eta S_a'(\psi)},$$

so implicit Euler is ...

$$\boldsymbol{\psi}^{n+1} = \boldsymbol{\psi}^n + h \frac{N(\boldsymbol{\psi}^{n+1})}{\boldsymbol{S}_s \boldsymbol{S}_a(\boldsymbol{\psi}^{n+1}) + \boldsymbol{\eta} \boldsymbol{S}_a'(\boldsymbol{\psi}^{n+1})},$$

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performs poorly:

• small denominator,

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- small denominator is squared for the Jacobian,

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- small denominator,
- small denominator is squared for the Jacobian,
- leading to many solver failures, which
- result in very small timesteps.

DAE formulation

$$S_s S_a(\psi^{n+1})(\psi^{n+1} - \psi^n) + \eta (S_a(\psi^{n+1}) - S_a(\psi^n)) = h N(\psi^{n+1}).$$

This is a lot better,

- larger time steps,
- happier nonlinear solver,
- error control easier to understand, and
- what most folks do.

ERDC ADH Code

- approximate VG-Mualem formulae with PL splines
- We explain the success of
 - finite difference approximation of Jacobians
 - Newton's method for implicit time-stepping
 - first order error estimation and control

ADH temporal integration

Solve

$$F(u) = S_s S_a(u)(u - \psi^n) + \eta \left(S_a(u) - S_a(\psi^n)\right) - hN(u) = 0,$$

with Newton's method.

ADH temporal integration

Solve

$$F(u) = S_s S_a(u)(u - \psi^n) + \eta \left(S_a(u) - S_a(\psi^n)\right) - hN(u) = 0,$$

with Newton's method. Approximate F'(u) by a finite difference Jacobian $\partial_h F(u)$

$$u_+ = u_c - (\partial_h F(u_c))^{-1} F(u_c),$$

and you get good results. Why?

Nonsmooth Calculus

 $F \in LIP$ implies F differentiable a.e. The generalized Jacobian (Clarke) at u is

$$\partial F(u) = \operatorname{co}\left\{\lim_{u_j \to u; u_j \in D_F} F'(u_j)\right\}$$

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$$u_{n+1} = u_n - V_n^{-1}F(u_n)$$

where $V_n \in \partial F(u_n)$. How do you compute V_n ? Can you use this stuff in the real world?

Piecewise smooth function: $\phi = \phi_l + \phi_r$ $\partial \phi(0) = [\phi'_l(0), \phi'_r(0)]$, a SET.


Difference approximations

Scalar functions

$$\partial_h \phi(u) = \frac{\phi(u+h) - \phi(u)}{h}$$

For Lipschitz functions:

$$\partial_h \phi(u) \in \partial \phi(\bar{u}) + O(h)$$

where $|u - \bar{u}| \le h$. Same story for scalar constitutive laws in PDEs. If you differentiate in coordinate directions!

Difference approximation accuracy $\phi'_l(0) + O(h) \le \partial_h \phi(u) \le \phi'_r(0) + O(h)$, so $\partial_h \phi(u) \in \partial \phi(0) + O(h)$



Semismoothness

A Lipschitz function F is semismooth (Mifflin, Pang, Qi) if

$$\lim_{w \to 0, V \in \partial F(u+w)} \frac{\|F(u+w) - F(u) - Vw\|}{\|w\|} = 0.$$

and semismooth of order 1 at *u* if

$$F(u+w) - F(u) - Vw = O(||w||^2)$$

for all $w \in R^N$ and $V \in \partial F(u+w)$ as $w \to 0$. What you need for local convergence of Newton's method. Piecewise smooth functions are semismooth of order 1.

Why semismoothness?

lf

- *F* semismooth of order 1,
- $F(u^*) = 0$, and
- everything in $\partial F(u^*)$ uniformly nonsingular,
- u_c near u^* ,

then if

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, where $V \in \partial F(u_c)$,

you get fast local convergence

$$||u_+ - u^*|| = O(||u_c - u^*||^2).$$

Convergence Proof, $e = u - u^*$

Semismoothness ($u \leftarrow u^*, w \leftarrow e_c, u + w \leftarrow u_c$) implies

$$F(u_c) - Ve_c = O(||e_c||^2)$$

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to get

$$e_{+} = e_{c} - V^{-1}F(u_{c}) = e_{c} - e_{c} + O(||e_{c}||^{2}) = O(||e_{c}||^{2}).$$

So what's up with ADH?

$$u_+ = u_c - (\partial_h F(u_c))^{-1} F(u_c)$$

and

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which implies

$$e_+ = O(||e_c||^2 + ||e_c||h+h).$$

Looks just like Newton if $||e_c|| >> \sqrt{h}$.

Iterative Linear Solvers

ADH uses preconditioned Krylov linear solvers. Termination on small relative linear residual,

$$\|F(u_c) + \partial_h F(u_c)s\| \leq \eta_c \|F(u_c)\|.$$

Convergence,

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Tradeoffs:

- Keep η small (accurate Newton step), for nonlinear performance,
- but not too small, to minimize linear solver cost.

Optimal difference increment

 ε_F : error in evaluation (eg floating point roundoff) Include this in V to get

 $V(u) \in \partial F(\bar{u}) + O(h + \frac{\varepsilon_F}{h})$

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So, if $||e_n|| = \sqrt{h}$, then

$$e_{n+1} = O((h + \frac{\varepsilon_F}{h}) ||e_n|| + ||e_n||^2 + h)$$

$$=O\left(\frac{\varepsilon_F}{h^{1/2}}+h\right)$$

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$$e_{n+1} = O((h + \varepsilon_F / h) ||e_n|| + ||e_n||^2 + h)$$

= $O\left(\frac{\varepsilon_F}{h^{1/2}} + h\right)$

which is minimized if $h = O(\epsilon_F^{2/3}) \approx 10^{-10}$ in IEEE.

Process: for u' = F(u), *F* Lipschitz continuous Goal: local truncation error $< \tau$.

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• Compare u^{n+1} and u^p to estimate error and change step size.

$$L = 2||u^{n+1} - u^p|| / |2h_n^2 - h_n h_{n-1}|$$

• Estimate Lipschitz constant of *u*' by

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$$h_n$$
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 - h_n ok? Enforce $Lh_{n+1}^2/2 < .9\tau$
- Completely rigorous if we're getting the Lipschitz constant right.

Numerical Experiments

Compare

$$L_{n+1} = 2\|u^{n+1} - u^p\|/|2h_n^2 - h_nh_{n-1}|$$

with

$$L(u^{n+1}) \frac{\|F(u^{n+1}) - F(u^n)\|}{t_{n+1} - t_n}$$

You want to see

$$r_n = \frac{L_n}{L(u^n)} \ge 1.$$

Study, RE for two media with 1 < n < 2.

Media Properies

Doromotor		ailt
Parameter	clay	SIII
n	1.09	1.37
α	0.244	0.478
S_r	0.179	0.074
η	0.33	0.40
K_s	1.10808e-5	1.1801e-03
T _{final}	600 days	150 days
maxh	10 days	5 days
$\tau - 10^{-2} h_0 - 10^{-9}$		

 $u = 10^{-1}, n_0 = 10^{-1}$

Clay



Silt



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Conclusions

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- Semismoothness is all you need for Newton's method.
 - Exotic math; software needn't know about it.
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- High-order methods in time seem to work. Why?